

From eq 1, $Pu = CRT$; hence

$$\lambda = \frac{P\varphi \sum x C_x}{C} = P\varphi Z_n \quad (22a)$$

or

$$Z_n = \frac{\lambda}{P\varphi} \quad (22b)$$

for the discontinuous case. This result is analogous to that in the continuous case where

$$Z_n = L/JP\varphi \quad (22c)$$

From eq 15 and 22b one can see that when (and if) $Z_n = Z_w$, then $\lambda = P/v$; but there is no necessity for this to occur at the pressure corresponding to the minimum in Z_n . In Table II we have summarized the relevant equations for the continuous and discontinuous cases.

It can be seen by examining Table II that the definition of Z_n has not changed in going from the continuous case to the discontinuous case, but that the definition of Z_w has changed due to the appearance of the new variable λ . The problem now has become the evaluation of the constants of integration. As yet this is not possible in all cases.

Table II: Summary of Equations for the Degree of Aggregation, Z_n and Z_w for the Continuous and Discontinuous Cases^a

| Continuous case | Discontinuous case |
|--|--|
| $Z_n = \frac{L}{JP\varphi}$ | $Z_n = \frac{\lambda}{P\varphi} = \frac{L_v}{J_v P\varphi}$ |
| $= \frac{1}{APve^{v/J}}$ | $= \frac{1}{A'Pve^{v/J}}$ |
| $= \frac{RTw/M^0}{EPve^{v/J}}$ | $= \frac{RTw/M^0}{E'Pve^{v/J}}$ |
| $Z_w = 1/v\varphi$ | $Z_w = 1/v\varphi$ |
| $= \frac{J}{ALv^2e^{v/J}}$ | $= \frac{1}{A'\lambda Pv^2e^{v/J}}$ |
| $= \frac{JRTw/M^0}{ELv^2e^{v/J}}$ | $= \frac{RTw/M^0}{E'\lambda v^2e^{v/J}}$ |
| $\varphi = -\left(\frac{d \ln C_1}{dv}\right)$ | $\varphi = -\left(\frac{\partial \ln C_1}{\partial v}\right)_\alpha$ |
| $= (AvLe^{v/J})/J$ | $= A'\lambda ve^{v/J}$ |

^a The constant A is different for the continuous and discontinuous cases. J and L are experimental and independent of which case is chosen for analysis. ^c $\lambda = L_v/J_v$ and it is a variable with pressure.

Further work is in progress on the mathematical and conceptual development of this theory.